

Investigation of the physical properties and Mulliken charge distribution of the cube perovskite BiGaO₃ is calculated by GGA-PBE

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We systematically studied the structure, electronic, elastic and optical properties of BiGaO₃ type perovskite cubes. We also report the calculation of all properties listed in the title of GGA-PBE approximation to BiGaO₃. This material has a band gap which can be considered as an indirect band gap. The maximum valence band is represented by the M point (VBM), whereas the conduction band is represented by the X point (CBM). The electronic structure of BiGaO₃ shows that it has a semiconductor indirect band gap of 1.37 eV. The elastic constant was determined in equilibrium confirming its accuracy. The bulk, shear, and Young's modulus were all extracted from the data. Poisson's ratio was also found. The confirmed structural parameters were in good agreement with the previously calculated experimental data. Our research suggests that BiGaO₃ is a promising piezoelectric, multiferroic, ferroelectric, and photo catalytic material.

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1. Introduction

Various composite oxides were extensively studied. However, perovskite oxide is especially versatile, showing an extent of different applications. Studies on this topic have been ongoing for quite some time now [1]. Recent studies have shown that perovskites, which include materials with Bi-metal transitions, have a lot of potential for various applications. [2]. Perovskite oxides are gaining more attention due to their technical applications such as electronic devices, light waves, piezoelectricity and iron [3-5]. A group of materials commonly used for lead-based iron and piezoelectric devices (such as PbTiO₃ and PbZrO₃). Due to the toxicity of lead, it was necessary to find materials that have similar properties and do not pollute the environment. Several attempts have been made in recent years for a solid BiGaO₃-based solution. Bi-based oxide compounds containing the non-magnetic element Ga have been receiving much attention due to the materials solid solution can enhance the photoelectric features of lead titanate and reduce the lead content [6]. Bismuth gallate BiGaO₃ was produced experimentally using high-pressure and high-temperature technologies [7]. Investigations on the physical properties of bulk and nanostructures of these compounds are limited, when compared to other multi-ferroic materials. Theoretically, the piezoelectricity and large ferroelectric polarization of BiGaO₃ were predicted [4,8,9]. Using the basic principles of density functional theory calculations to determine the electronic properties, structure and instability of the four cubic perovskites BiMO₃ with M (M = Ga, Al, In and Sc). A research has shown that BiMO₃ oxides or modified versions are promising photo catalysts, ferroelectric materials, piezoelectric materials, and multiferroic materials [10,11]. In this study, we focused on the physical properties of BiGaO₃ using density functional theory. This is important because it allows us to understand the compound completely, which is necessary for its use in technological applications. Elastic constants provide valuable information about the anisotropic character of the bonding and structural stability.

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2. Method of calculations

It is assumed that the considered BiGaO_3 has an ideal cubic perovskite structure with space group $\text{Pm}\bar{3}\text{m}$ (#221). Here, the atomic positions of the primitive units are represented by Bi:1a(0,0,0), Ga:1b(0.5,0.5,0.5) and O:3d(0.5,0,0) are defined. The electronic configuration of the compound is assumed to be $4f^{14}5d^{10}6s^26p^3$, $3d^{10}4s^24p^1$, $2s^22p^2$ for Bi, Ga and O respectively. The Kohn-Sham density functional theory is a well-known and reliable method for simulating the ground state characteristics of atoms, molecules, and solids. This theory can be accurate without knowing the exact form of the electron density; because many-body problems are complex and nonlinear, approximations are made to make the theory tractable[12]. A generalized gradient approximation is one of the most popular functional approximations (GGA)[13]. The GGA are now ready, and they contain PBE[14]. The Broyden–Fletcher–Goldfarb–Shanno minimization method was used for structural optimization (BFGS), the force on each ion converged to less than $0.01\text{eV}/\text{\AA}$. A cut-off of the kinetic energy of (800) eV and a total energy convergence threshold of (1×10^{-6}) eV were used. The Brillouin zone integrations were performed by using a Γ -centered ($12 \times 12 \times 12$) k-point grid.

3. Results and discussion

3.1. Geometry optimization and Band structure

The BiGaO_3 compound crystallizes in the $\text{Pm}\bar{3}\text{m}$ (#221) space group, which is ideal for cubic perovskites. (5 atoms) per unit cell. Figure (1) shows the crystal structure of BiGaO_3 at room temperature. It can be noticed that Bi cation appears at the corner of the cube, with the oxygen ion in the center of the plane in the center cation. This indicates that BiGaO_3 has achieved a cubic structure. The structural parameters of the BiGaO_3 cube have been calculated. The results are summarized in Table (1) which shows that the lattice constant (a) is overestimated by about 0.5%. These deviations are within the permissible range of first-principles error. Our findings are similar to these of other researchers who have used the full potential linearized augmented plane wave (FP-LAPW) method[10,11]. Table (1) includes these values as well as other theoretical results.

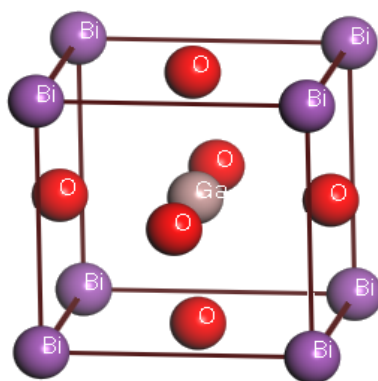


Fig. 1. Cube lattice structure of BiGaO_3 .

Table 1. Calculated equilibrium lattice parameters, volume and band gap for BiGaO_3 compound in comparison with the available theoretical results.

Parameter	$a(\text{\AA})$	$V(\text{\AA})^3$	E_g (eV)
this work	3.8573	57.389	1.37
[10]	3.905	59.547	1.34
[11]	3.816	55.568	1.41

Figure (2) describes the energy band gap of BiGaO₃. The GGA-PBE technique is a powerful way to determine the band structure along the high symmetry Γ , R, M direction of, and the Brillouin zone represented in X. BiGaO₃ compound has an indirect energy band gap (1.37 eV). Figure (2) Shows that the CBM and VBM are located at points that are completely different from each other. points M ($k = 2\pi / a (110)$), and X in the Brillouin zone ($k = 2\pi / a (100)$). Our indirect band gap (MX) calculation is consistent with the values of the other calculations listed in Table (1).

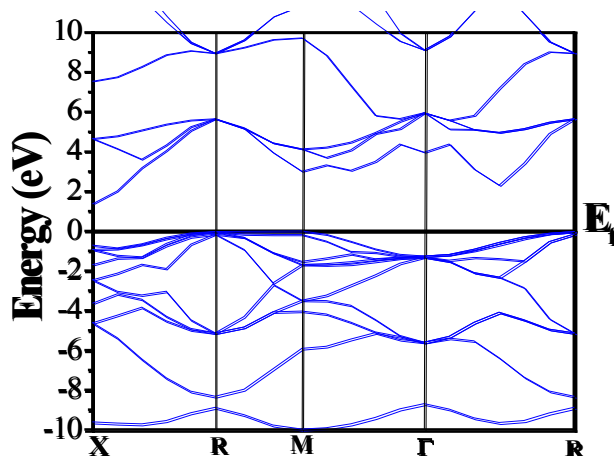


Fig. 2. The energy band structure of a BiGaO₃ compound has been calculated.

3.2. Density of states (DOS)

To clarify the nature of the electronic band structures of BiGaO₃, the total density (TDOS) and partial density of states (PDOS) were calculated. It is shown in figure (3) that PDOS can determine the nature of the angular momentum of various structures. From figure (3), the valence bands VB can be divided into four different energetic regions. The first region of the energy spectrum extending from -22.8 to -22 eV which includes Bi-5d. The second region starts from an energy range of -8.5 to -18.7 eV, which is dominated by the O-2 state that has a small contribution to the state of Ga-3d. A third region from -14.7 to -12.4 eV, consisted of Ga-3d orbital states. Finally, it is clear that dominating states of O-2p are in the range of 0 and -6.6 eV (upper VB), with a mixture of Ga-sp and Bi-6P orbitals. The hybridization of Ga atoms and O atoms in this region is located near the Fermi level indicates that the Ga-O bond is a covalent bond. The lower part of the band gap of BiGaO₃ is mostly filled by the 6p state of Bi, while the upper part is mostly filled by the s and p states of Ga. The s and p electrons of Ga force the Bi-6p state closer to the Fermi level. Our calculations of the band structure, TDOS and PDOS of BiGaO₃ are consistent with previous theoretical results [10,11].

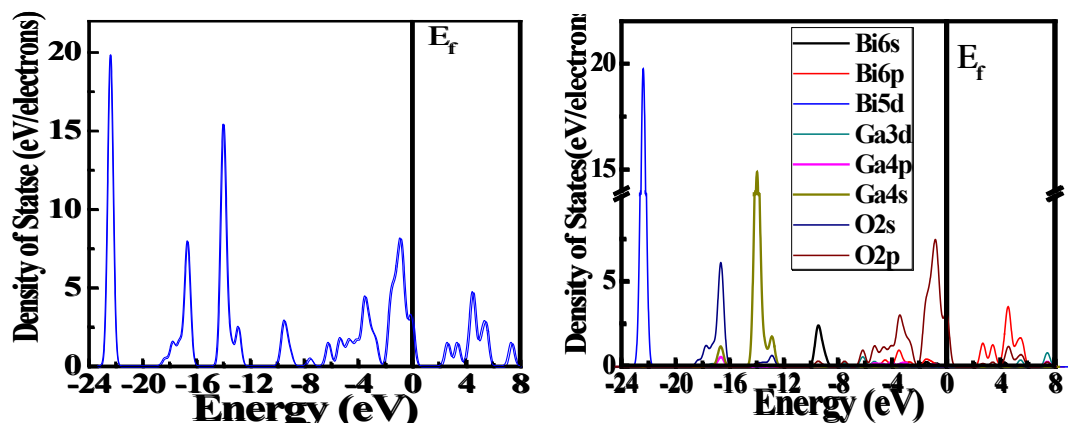


Fig. 3. Calculated total density of states (TDOS) and partial density of states (PDOS) of BiGaO_3 at the predicted equilibrium lattice constant.

We have also used the Mulliken charge collection of BiGaO_3 to describe the behavior of bonds. Table (2) shows the Mulliken population charge values. Due to the transferred charge from Bi to O is approximately 1.51 electrons and the transfer of the charge from Ga to O is 1.29 electrons, we deduced that BiGaO_3 bonding characteristic is a mix of ionic and covalent nature. Furthermore, the bond between Ga-O is a covalent bond that is highly strong, which stronger than the bond between Bi-O.

Table 2. Mulliken charge population of cubic perovskite BiGaO_3 .

Species	s	p	d	f	Total	Charge (electron)
O(1)	1.86	5.07	0.00	0.00	6.93	-0.93
O(2)	1.86	5.07	0.00	0.00	6.93	-0.93
O(3)	1.86	5.07	0.00	0.00	6.93	-0.93
Ga(1)	0.61	1.11	10.00	0.00	11.71	1.29
Bi(1)	2.00	1.48	0.00	0.00	3.45	1.51

3.3. Optical properties

Figure (4) depicts the estimated optical characteristics at the equilibrium lattice constant for energy ranges up to 30 eV. In order to calculate the observed structures in the spectrum, one must consider the transition between occupied and unoccupied energy bands in the structure of the electronic energy bands at high symmetry points in the Brillouin zone and must be considered According to the explanation[15]. The dielectric properties are expressed by using the Kramer–Krung relationship for the complex dielectric function $\epsilon(\omega) = \epsilon_1(\omega) + i\epsilon_2(\omega)$ [16]. The first part, $\epsilon_1(\omega)$ is the real part, which determines the polarization and scattering of light, and the second part, $\epsilon_2(\omega)$ is the imaginary part that describes the absorbance of the material. The real part of dielectric function $\epsilon_1(\omega)$ is shown in Figure (4(a)) in the energy range (0-15)eV. The static value of $\epsilon_1(0)$ is (5.7) is observed for BiGaO_3 , $\epsilon_1(\omega)$ continues to increase from $\epsilon_1(0)$, reaching the value at maximum level before decreasing below zero at specific ranges of energy. The incident photon beam at these energy bands is completely weakened where a negative value in this region means a loss in energy transmission as well as a weakening in light transmission. Based on the $\epsilon_2(\omega)$ spectra in Figure 4(a), it is clear that the first critical points for BiGaO_3 occur at about 2.75eV. This agrees with our predictions. The fundamental absorption edges [17] are caused by the direct inter-band transitions VBM and CBM in the BiGaO_3 compound. Electrons travelling from O at the 2p states in the valence band region to the Ga at the 4s states in the conduction band

region after (2.9 and 4.6)eV. Our calculated optical and electronic properties suggest that the investigated compounds are possible to be implemented in photovoltaic cell fabrication, according to practical studies[18].The extinction index and refractive index are found to help analyze the degree of transparency of the studied material and the ability of absorbing incident light (photons).The relation of extinction coefficient $k(\omega)$ and refractive index $n(\omega)$ with real part are as follows:

$$k(\omega) = \sqrt{\frac{|\varepsilon(\omega)| - \varepsilon_1(\omega)}{2}} \quad (1)$$

$$n(\omega) = \sqrt{\frac{|\varepsilon(\omega)| + \varepsilon_1(\omega)}{2}} \quad (2)$$

Figure 4b shows the refractive index $n(\omega)$ and extinction coefficient $k(\omega)$ of BiGaO₃. Over the wide-range energy over 15eV, the values of refractive index spectrum of $n(\omega)$ is very broad. $n(\omega)$ and $\varepsilon_1(\omega)$ seem to have the same character. It can be shown in Figure 4b that $n(\omega)$ increases as the energy peaks at 4.45 eV for BiGaO₃.The ultraviolet is represented by a peak in the transparency region, which then descends to a minimum level at 7.5eV. BiGaO₃'s static refractive index $n(0)$ varies inversely with the band gap, with calculated values of approximately 2.44.Fig. (4d) shows that extinction coefficient $k(\omega)$ local peaks correspond the static component of $\varepsilon_1(0)$.The final behavior of the $k(\omega)$ and $n(\omega)$ peaks is determined to be similar for the compound and in good agreement with another anticipated investigation of the $k(\omega)$ and $n(\omega)$ peaks [19].The wavelength-dependent reflectance coefficient $R(\omega)$ for BiGaO₃ is shown in Figure (4c). $R(0)$ is the static reflectivity. The first peak is at 3.4eV.It is worth noting that the reflectivity achieves a maximum when $\varepsilon_1(0)$ is smaller than zero. The material shows the metallic characteristic for negative values of $\varepsilon_1(\omega)$ [20].Because a compound's reflectivity is proportional to its metallicity, its value $R(0)$ is greatest for negative $\varepsilon_1(\omega)$ As shown below [21], the absorption coefficients $\alpha(\omega)$ reflect the influence and efficiency of any material in absorbing energy packets (photons) with respect to photons of light of energy $E=h\nu$.

$$\alpha(\omega) = \sqrt{2}\omega \left[\sqrt{\varepsilon_1^2(\omega) - \varepsilon_2^2(\omega) - \varepsilon_1(\omega)} \right]^{\frac{1}{2}} \quad (3)$$

The BiGaO₃ absorption coefficients $\alpha(\omega)$ of are shown in Fig (4d). The significant absorption zones of the examined materials are supported by comparing with Figure (4a) $\varepsilon_1(\omega)$, BiGaO₃ should be optically active in these locations. According to the computed data, the absorption coefficient for BiGaO₃ exhibits a sharp edge about 16 eV. These sharp absorption coefficients edges are similar to the band gap of semi-conductive materials, which allow only photons with sufficient energy to stimulate an electron to the unoccupied levels of the conduction band from the occupied levels of the valence band to be seen into the material. The presence of high absorption zones of 12-17eV perovskite, may indicate that they are promising candidates for optoelectronic device applications in the UV range. The observed optical conductivity Fig. (5e), Applying energy will result in a response in optical conduction after 2.35eV for BiGaO₃ respectively. The BiGaO₃ optimum optical conductivity is approximately at 5.6eV. The $L(\omega)$ spectra of BiGaO₃ are shown in figure (4f), it is obtained that the maximum peak is located at 22 eV. The electron energy loss function $L(\omega)$ plays an important role in describing the energy lost for a fast electron passing through a material. The prominent peaks in the $L(\omega)$ spectra represent the property associated with plasma resonance (mass oscillation of valence electrons) and the corresponding frequency is the so-called plasma frequency $\omega(p)$ [22].The peaks of $L(\omega)$ correspond to the trailing edges in the reflection spectra, where the occurrence of prominent peaks of $L(\omega)$ figure (4f) can be observed at energies corresponding to the abrupt reductions of $R(\omega)$ figure (4c).The function for calculating electron energy loss $L(\omega)$ is used to describe the energy lost by a rapid electron traveling through a materials. The large peaks in the $L(\omega)$ spectra

show the plasma resonance property (mass oscillation of valence electrons), and the related frequency is known as plasma frequency $\omega(p)$ [22].

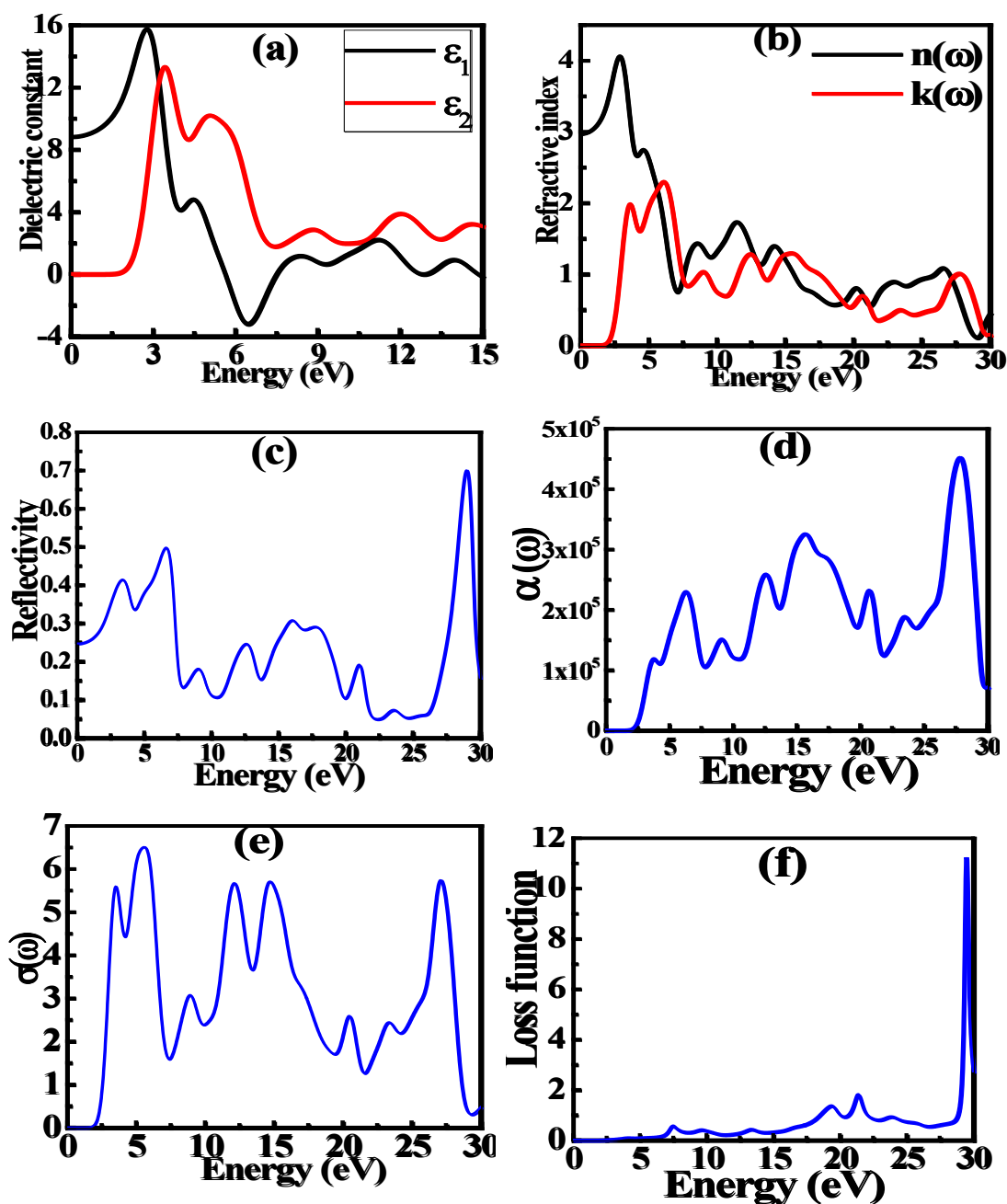


Fig. 4. Calculated optical constants for BiGaO₃: dielectric function: (a) real part $\epsilon_1(\omega)$ and imaginary part $\epsilon_2(\omega)$ (b) The extinction index $K(\omega)$ and the refractive index $n(\omega)$, (c) The reflectivity $R(\omega)$, (d) the absorption coefficients $\alpha(\omega)$, (e) The photoconductivity $\sigma(\omega)$, and (f) loss function $L(\omega)$.

3.4. Elastic properties

The elastic constants, which play a key role in supplying valuable information on the correlation property between neighboring atomic levels, determine the qualities of a material that endures deformation and stresses and restores its original shape after the stress is removed. In general, the importance of the elastic constants of solids lies in the fact that they provide a link between the mechanical and dynamic behavior of crystals. Table 3 shows the elastic moduli of BiGaO₃ calculated using the perturbation method at equilibrium pressure within the GGA-PBE

approximation . In this table we present the calculated values of the elastic constants C_{ij} . The cubic system is characterized by three different symmetry elements (C_{11} , C_{22} , and C_{44}), in cubic crystals, the elastic contrast factor (A) is a measure of the solids elastic anisotropy degree, defined as $A = 2C_{44}/(C_{11} - C_{12})$, the material is completely isotropic If the value of $A = 1$. If A value of is less or greater than 1 it means that the value of elastic anisotropy degree is shown. Our calculations showed Zener anisotropy for BiGaO₃ is less than 1, indicating that BiGaO₃ is an anisotropic material, according to what the researchers also explained [23]. Moreover, Bulk modulus B is a linear combination of elastic constants (C_{11} and C_{12}) for cubic systems. Where C_{ij} is calculated as a function of the volume-conserving stress that breaks the symmetry. Reference [24] provides more information about the accounts. Using the following expression, the internal strain parameter (ζ) can be calculated:

$$\zeta = \frac{C_{11} + 8C_{12}}{7C_{11} + 2C_{12}} \quad (4)$$

The upper limit of (ζ) corresponds to the insignificant contribution of the bond stretching/contracting to resist external stress., while the lower limit represents the insignificant contribution of the bond bending to resist the external stress. For crystalline solids of cubic shape, there are three types of independent elastic constants (C_{11} , C_{22} and C_{44}) that are used to check the mechanical stability according to Born criteria such as given in equation (5).

$$C_{11} - C_{12} > 0; C_{11} + 2C_{12} > 0; C_{11} > 0; \text{and } C_{44} > 0 \quad (5)$$

Table (3) The computed elastic constants for BiGaO₃ display that the unidirectional elastic consistent C_{11} , that's associated with unidirectional strain alongside the main crystal direction, is a good deal better than C_{44} , which suggests that BiGaO₃ has a weaker resistance to pure shear deformation as compared to the unidirectional compressive resistance. In a cubic crystal, there are conditions for the classical mechanical stability of elastic constants. The estimated elastic constants in table (3) satisfied the above-mentioned stability conditions, suggesting that BiGaO₃ is mechanically stable. After evaluating the elastic constants (C_{11} , C_{22} , and C_{44}) for BiGaO₃, we can estimate relevant properties such as shear modulus (G), bulk modulus B, Poisson's ratio (ν), Young's modulus (E), Lamé's coefficients (μ, λ), and Pugh ratio (G/B) [25-27] by using the following relations:

$$G = \frac{G_V + G_R}{2} \quad (6)$$

$$C_V = \frac{C_{11} - C_{12} + 3C_{44}}{5} \quad (7)$$

$$C_R = \frac{5(C_{11} - C_{12})C_{44}}{4C_{44} + 3(C_{11} - C_{12})} \quad (8)$$

$$\nu = \frac{3B+E}{6B} \quad (9)$$

$$E = \frac{9BG}{3B+G} \quad (10)$$

$$\mu = \frac{E}{2(1+\nu)} \quad (11)$$

$$\lambda = \frac{\nu E}{(1+\nu)(1-2\nu)} \quad (12)$$

Table (3) corresponds to the present value of the shear modulus and for the bulk modulus of BiGaO₃ (B = 205.72 GPa and G = 104.32 GPa). Teter [28] illustrated that high values of shear modulus and bulk modulus are a guide and indicator of hardness. It is worthy to note that covalent and ionic bonds are distributed between cations that might be octahedral or cube-octahedral rounded by O atoms in BiGaO₃ compound [29]. The calculated results reveal that the nature BiGaO₃ compound is ionic, as the Poisson ratio calculated to be close to 0.3, and it was found that the calculated G/B value for BiGaO₃ is equal to (0.5), which also shows this compound bonding is of an ionic nature, where it can reveal (G/B) On the brittle or ductile behavior of the studied compound, it was shown from the ratio (G/B) that it has a strong ductile character [30,31]. The strong ductile character and ionic properties observed for BiGaO₃ compound indicate its suitability for practical device applications.

Table3. Calculated elastic parameters C_{ij} (GPa), Young's modulus E (GPa),shear modulus G (GPa),bulk modulus B (GPa),Pugh'sindex G/B , Poisson's ratio(ν) internal strain parameter (ζ), anisotropy factor A ,and Lamé's coefficients (μ,λ) ofBiGaO₃.

parameter	C_{11}	C_{12}	C_{44}	E	G	B	G/B	ν	ζ	A	λ	μ
this work	364	126	95	269	105	205	0.5	0.28	0.49	0.8	134	105
Ref[32]	387.7	128.9	119	309.9	123	215.1	0.57	0.74	0.48	0.92	274.5	89
Ref[33]	375.8	129.7	98.3	267.7	104.3	205.7	0.51	0.28	0.5	0.86	136.7	103.3

4. Conclusion

The GGA-PBE approximation was used to calculate the optical properties, electronic structures, elastic constants, and the structural parameters of BiGaO₃. The computed lattice constants of BiGaO₃ agree with other theoretical results. The calculated crystal structure of BiGaO₃ indicates that it possesses a cubic crystal structure with a space group pm-3m(221#). Moreover, the electronic structures revealed that the upper part of the valence bands were identified by the O-2p states, and the lower part of the conduction bands were identified by the Bi-6p states. Electronic structure calculations showed that BiGaO₃ is a semiconductor with an indirect M-X band gap separating the conduction bands (CB) from the valence bands (VB). The value of the calculated elastic constants for C_{ij} is positive, indicating that BiGaO₃ is characterized by mechanical stability.

The current results assume the character of ionic bonding of the studied compound and lead to be ductile material classified. The real and imaginary parts of the dielectric functions, extinction coefficient, refractive index, reflectivity spectra, electronic energy loss function, optical absorption coefficient, real part of the optical and optical conductivity of the composite were determined. It is observed that $n(\omega)$ and $k(\omega)$ have almost the same properties as $\epsilon_1(\omega)$ and $\epsilon_2(\omega)$, respectively, BiGaO₃ has interesting as piezoelectric, ferroelectric, multiferroic, and photo catalytic materials.

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