

## EFFECT OF Bi ADDITIVE ON ELECTRICAL PROPERTIES OF CHALCOGENIDE $\text{Se}_{80}\text{Te}_{20}$ AMORPHOUS THIN FILMS

K. SHARMA<sup>a</sup>, M. LAL<sup>b</sup>, A. KUMAR<sup>a</sup>, N. GOYAL<sup>a\*</sup>

<sup>a</sup>*Department of Physics, Panjab University, Chandigarh, India-160014.*

<sup>b</sup>*GGSDS College, Chandigarh, India-160030*

The electrical properties of  $\text{Se}_{80-x}\text{Te}_{20}\text{Bi}_x$  ( $x = 0, 2, 4, 6, 8$ ) amorphous thin films prepared by thermal evaporation technique have been studied. The dark conductivity ( $\sigma_d$ ), charge carrier concentration ( $n_\sigma$ ) and photoconductivity ( $\sigma_{ph}$ ) increases as the concentration of Bi additive increases upto  $x = 4$  at. % in  $\text{a-}\text{Se}_{80-x}\text{Te}_{20}\text{Bi}_x$  thin films. The activation energy ( $\Delta E$ ) and photosensitivity ( $\sigma_{ph}/\sigma_d$ ) decreases sharply after Bi incorporation. As the Bi concentration further increases (when  $x = 6$  at. %) a reverse in the trend has been observed for all these parameters. With further increase of Bi content in  $\text{Se}_{80-x}\text{Te}_{20}\text{Bi}_x$  amorphous thin films (when  $x = 8$  at. %), the dark conductivity ( $\sigma_d$ ), charge carrier concentration ( $n_\sigma$ ) and photoconductivity ( $\sigma_{ph}$ ) again increases and activation energy ( $\Delta E_d$ ) and photosensitivity ( $\sigma_{ph}/\sigma_d$ ) decreases. Intensity dependence of photoconductivity has also been studied at 303 K and follows a power law as  $\sigma_{ph} \propto F^\gamma$ , where  $\gamma$  lies between 0.5 and 1. The differential life time ( $\tau_d$ ) is determined from the decay of photocurrent with time. The differential life time increases as the Bi concentration increases upto  $x = 4$  at. % and decreases with further increase of Bi content (when  $x = 6$  at. %) and then again increases as Bi content in  $\text{Se}_{80-x}\text{Te}_{20}\text{Bi}_x$  amorphous thin films further increases upto  $x = 8$  at. %. The results are explained on the basis of increase in the density of localized states present in the mobility gap of  $\text{Se}_{80-x}\text{Te}_{20}\text{Bi}_x$  thin films.

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### 1. Introduction

Chalcogenide glasses are a very important class of semiconductor materials because of their technological applications in xerography, switching and memory devices [1-4]. These glasses have attracted much attention for the development of active and passive infrared devices. These glasses are low-phonon energy materials and are generally transparent from the visible to infrared region [5-7]. Chalcogenide glasses are known to have flexible structure, in the sense that each atom can adjust its neighboring environment to satisfy the valance requirements. Chalcogenide glasses are sensitive to the absorption of electromagnetic radiation and show a variety of photo-induced effects and various models have been put forward to explain these effects, which can be used to fabricate wave-guide and fiber structures [8-12]. As chalcogenide glasses have poor thermo-mechanical properties, in order to enlarge their domain of applications, it is necessary to increase their softening temperature and mechanical strength. The interest in these materials arises particularly due to their ease of fabrication in the form of bulk and thin films. A typical chalcogenide has a relatively sharp optical absorption edge, single electrical activation energy, efficient photo excited conductivity and luminescence. The photoconductivity of these glasses is helpful in understanding the phenomenon of recombination kinetics, and the nature and distribution of localized states in the forbidden energy gap. In chalcogenide glasses, the localized states in the mobility gap are the  $D^+$  and  $D^-$  states [13]. Transient photoconductivity measurements

\* Corresponding author: ngoyal@pu.ac.in

in chalcogenide glasses are important as the photocurrent rise and decay with time depends upon the presence of traps in the mobility gap of these materials. These traps originate from the defect states present in these materials. Therefore, such measurements made at different temperatures, intensities and illumination times give important information about the defect states [14].

In pure state Se has several disadvantages because of its short life time and low sensitivity. To overcome these difficulties, certain additives are used, e.g. Te, Sb, Bi, In, Ge, etc and binary and ternary alloys are formed. Recently, it has been pointed out that Se-Te alloys have more advantages than a-Se from the technological point of view due to their greater hardness, higher crystallization temperature, higher photosensitivity and smaller ageing effects [16-19].

The effect of an impurity in an amorphous semiconductor may be widely different, depending upon the conduction mechanism and the structure of material [21]. Metallic impurities such as Pb and Bi, produces a remarkable change in the electrical properties of chalcogenides glasses. The conductivity type switches from p to n type on the addition of Bi [22-24]. The third element behaves as chemical modifier as it is reported to expand the glass forming region. The addition of third element like Bi in Se-Te binary alloy is expected to change the electrical properties of host alloy.

In the present work, the photoconductive properties of thin films of the  $\text{Se}_{80-x}\text{Te}_{20}\text{Bi}_x$  system have been studied in detail. Temperature-dependent dark conductivity, photoconductivity and intensity dependent steady-state photoconductivity have been measured for five samples with different compositions. The decay of photocurrent has also been measured at room temperature for all compositions.

## 2. Experimental procedure

Glassy alloys of  $\text{Se}_{80-x}\text{Te}_{20}\text{Bi}_x$  ( $x = 0, 2, 4, 6, 8$ ) were prepared by melt-quench technique. High purity 99.999% Se, Bi and Te granules were weighted according to the formula of  $\text{Se}_{80-x}\text{Te}_{20}\text{Bi}_x$  ( $x = 0, 2, 4, 6, 8$ ). The powder mixture was loaded into quartz ampoule and sealed under vacuum at  $10^{-4}$  Pa. The sealed quartz ampoule was loaded in a furnace and heated to  $950^\circ\text{C}$  at a rate of  $3\text{-}4^\circ\text{C}/\text{minute}$  for 18 hours to ensure the composition homogeneity and quenched in liquid nitrogen. The ingots were crushed, separated and grounded.

Thin films were synthesized by thermal evaporation technique under high vacuum conditions ( $\sim 10^{-4}$  Pa) using a small piece of bulk alloy as a source material and glass as a substrate. The films were kept inside the deposition chamber for 24 hours to achieve the metastable equilibrium. The glassy nature of thin films was checked with the help of X-ray diffraction technique using  $\text{Cu-K}_\alpha$  radiation. For electrical measurements vacuum – evaporated indium electrodes were used. A coplanar structure (length  $\sim 1.8$  cm and separation between electrodes  $\sim 0.7$  mm) was used for present measurements.

A three terminal sample holder in which light could be shown through a transparent window has been used for the measurement of photoconductivity of thin films. A copper-constantan thermocouple has been inserted inside the sample holder and kept close to the sample to measure the correct temperature. A vacuum of the order of  $10^{-4}$  to  $10^{-5}$  Torr was achieved inside the sample holder using vacuum pump. The light source for these measurements was a 200 W tungsten lamp. Light intensity is measured by a digital Luxmeter (Testron, model TES-1332). The photocurrent ( $I_{\text{ph}}$ ) is obtained after subtracting dark current ( $I_{\text{d}}$ ) from the current measured in the presence of light. Before measurement, the films were first annealed below their glass transition temperature for two hours in a vacuum  $\sim 10^{-4}$  Torr for different samples and the dark conductivity, photoconductivity and transient photoconductivity measurements are carried out. The present measurements were made by applying only 20 V across the films using dc power supply (CROWN DC-regulated Power Supply 0-30V/2A) and the resulting current was measured with digital piccoammeter (Digital Piccoammeter DPM-111, Scientific equipment Roorkee).

### 3. Result and discussion

Fig. 1 shows the temperature dependence of dark conductivity ( $\sigma_d$ ) for a  $\text{Se}_{80-x}\text{Te}_{20}\text{Bi}_x$  ( $x = 0, 2, 4, 6$  and  $8$ ) thin films. The  $\ln\sigma_d$  versus  $1000/T$  curves are straight lines in the investigated temperature range 273K to 343K. This indicates that the conduction in these glasses is through an activated process having single activation energy in the investigated temperature range. The  $\sigma_d$  can, therefore, be expressed by the usual relation

$$\sigma_d = \sigma_0 \exp(-\Delta E/k_B T) \quad (1)$$

where  $\sigma_0$  is the pre-exponential factor,  $\Delta E$  is the activation energy for conduction and  $k_B$  is the Boltzmann constant. The values of  $\sigma_d$  at 303K and  $\Delta E_d$  have been calculated from the curves plotted in Fig. 1. The value of  $\sigma_d$  increases from  $1.01 \times 10^{-9} (\Omega\text{-cm})^{-1}$  to  $1.26 \times 10^{-7} (\Omega\text{-cm})^{-1}$  as the Bi concentration increases from  $x = 0$  to  $x = 4$  followed by a decrease to  $6.93 \times 10^{-8} (\Omega\text{-cm})^{-1}$  on increasing the Bi concentration to  $x = 6$  and then  $\sigma_d$  again increases to  $3.53 \times 10^{-7} (\Omega\text{-cm})^{-1}$  with further increase of Bi content to  $x = 8$  at 303K. Correspondingly, the value of  $\Delta E_d$  decreases from 0.55 eV to 0.46 eV when  $x$  increases from 0 to 4 and then  $\Delta E_d$  increases to 0.48 eV as the Bi composition further increases to  $x = 6$  and then with further increase of Bi content to  $x = 8$ ,  $\Delta E_d$  again decreases to 0.44 eV.

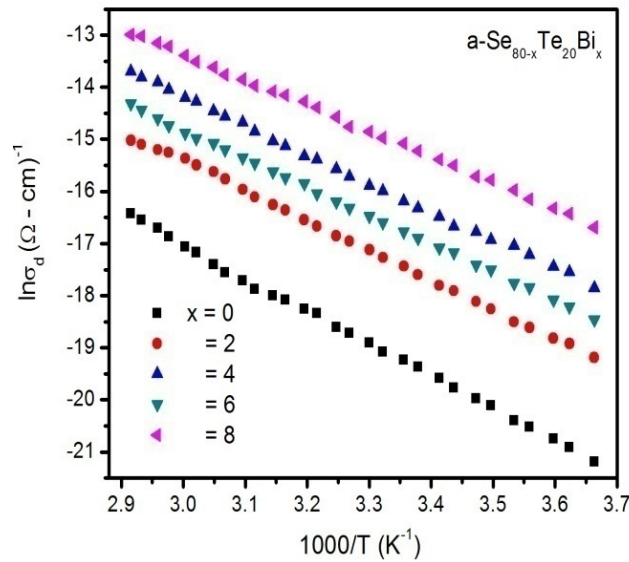


Fig. 1. Variation of  $\sigma_d$  with temperature for  $a\text{-Se}_{80-x}\text{Te}_{20}\text{Bi}_x$  thin films.

Fig. 2 shows the temperature dependence of steady state photoconductivity ( $\sigma_{ph}$ ) at a particular intensity (1043 lux) for all compositions. The value of  $\sigma_{ph}$  increases with temperature in the measured temperature range (273K to 343K). The activation energy of photoconduction is smaller than the dark conduction. The values of  $\sigma_{ph}$  have been calculated at 303 K and at intensity 1043 lux from the plots of Fig. 2 and given in Table 1. The values of  $\sigma_d$  and  $\sigma_{ph}$  follow the same pattern on Bi incorporation in the  $\text{Se}_{80}\text{Te}_{20}$  as shown in Table 1.

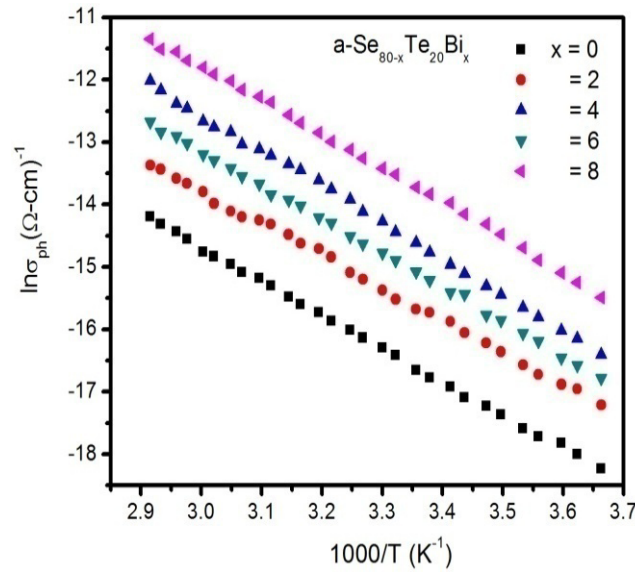


Fig. 2. Variation of  $\sigma_{ph}$  with temperature for  $a\text{-Se}_{80-x}\text{Te}_{20}\text{Bi}_x$  thin films.

The charge carrier concentration ( $n_\sigma$ ) for different samples can be calculated using the equation [25-26]

$$n_\sigma = 2(2\pi mk_B T/h^2)^{3/2} \exp(-\Delta E_d/k_B T) \quad (2)$$

where 'm' is the mass of the charge carrier and 'k<sub>B</sub>' is the Boltzmann constant. The calculated values of  $n_\sigma$  are given in the Table 1. It is evident from Table 1, that the values of  $n_\sigma$  increases with the increase of Bi concentration upto  $x = 4$  and then decreases as Bi content increases to  $x = 6$  and it again increases on increase in the concentration of Bi additive to  $x = 8$  at. %.

An important and useful parameter in the photoconductivity measurements is the photosensitivity ( $\sigma_{ph}/\sigma_d$ ), which decides the use of that material in the photoconductive devices. From the measured values of  $\sigma_{ph}$  and  $\sigma_d$ , we have therefore calculated the  $\sigma_{ph}/\sigma_d$  for all the samples at a particular intensity (1043 Lux) and at temperature 303 K. The  $\sigma_{ph}/\sigma_d$  decreases quite appreciably after the incorporation of Bi additive upto  $x = 4$ . As Bi content further increases ( $x = 6$ ), the value of  $\sigma_{ph}/\sigma_d$  increases to 6.07 followed by a decrease to 4.19 on increasing the Bi concentration to  $x = 8$  at. %.

An increase in dark conductivity with a corresponding decrease in activation energy is generally found to be associated with the shift of Fermi level in impurity doped chalcogenide glasses [27]. On the other hand, an increase in dark conductivity after impurity doping has also been explained in terms of the increased hopping conduction in impurity induced states [28]. The dark activation energy ( $\Delta E_d$ ) alone does not provide any indication as to whether conduction takes place in extended states above the mobility edge or by hopping in the localized states. So, the increase in  $\sigma_d$  after Bi incorporation can be interpreted in terms of either of the two above mentioned processes. According to the second interpretation, the density of localized states should increase which should lower the photosensitivity at a particular temperature and intensity, but in the first interpretation, the density of states should not increase. Since the photosensitivity ( $\sigma_{ph}/\sigma_d$ ) decreases sharply after the Bi incorporation indicating an increase in density of localized state, therefore, the first interpretation seems unlikely in the present case.

Table 1. Variation of electrical parameters at different Bi concentration in *a*-Se<sub>80-x</sub>Te<sub>20</sub>Bi<sub>x</sub> thin films

x	$\sigma_d$ ( $\Omega^{-1}\text{cm}^{-1}$ )	$\sigma_{ph}$ ( $\Omega^{-1}\text{cm}^{-1}$ )	$\Delta E_d$ (eV)	$\sigma_{ph}/\sigma_d$	$n_\sigma$ ( $\text{cm}^{-3}$ )	$\tau_d$ (sec)
0	$1.01 \times 10^{-9}$	$8.35 \times 10^{-8}$	0.55	82.67	$1.83 \times 10^{10}$	32
2	$3.57 \times 10^{-8}$	$2.98 \times 10^{-7}$	0.52	8.34	$5.77 \times 10^{10}$	58
4	$1.26 \times 10^{-7}$	$6.33 \times 10^{-7}$	0.46	5.02	$5.73 \times 10^{11}$	83
6	$6.93 \times 10^{-8}$	$4.21 \times 10^{-7}$	0.48	6.07	$2.66 \times 10^{11}$	67
8	$3.53 \times 10^{-7}$	$1.48 \times 10^{-6}$	0.44	4.19	$1.23 \times 10^{12}$	104

The photosensitivity ( $\sigma_{ph}/\sigma_d$ ) depends on the lifetime of the excess charge carriers which in turn depends on the density of localized states in a particular material. The higher the density of states, the lower will be lifetime, as these defect states may act as recombination centers in presence of light. This is found to be due to the formation of intimate valence alternation pairs (IVAPs) under illumination. The defects are converted by bond switching reactions to random pairs of charged defects, known as light induce metastable defects (LIMDs) [29-31]. Such an increase in the density of defect states after incorporation of Bi additive can be explained by assuming that Bi atoms act as impurity centers in the mobility gap and induces structural changes in the network which may disturb the balance of charged defects and consequently change the electrical conduction. New trap states are created in the mobility gap due to addition of Bi and concentration of  $D^+$  states get reduced as compared to  $D^-$  states. This increase in conductivity may be attributed to current due to electrons only [33]. When Bi content increases to  $x = 6$  at. %, then addition of Bi may satisfy the bonding states of Se-Te [34] and density of localized states decreases which leads to decrease in conductivity of Se<sub>74</sub>Te<sub>20</sub>Bi<sub>6</sub> as compared to Se<sub>76</sub>Te<sub>20</sub>Bi<sub>4</sub>.

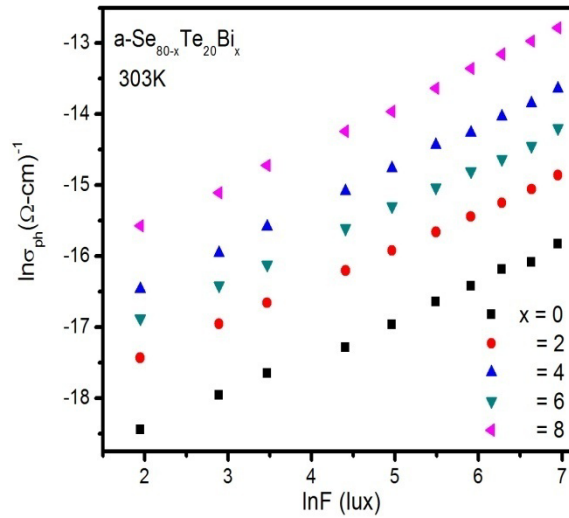


Fig. 3. Intensity dependence of photoconductivity for different *a*-Se<sub>80-x</sub>Te<sub>20</sub>Bi<sub>x</sub> glassy alloy thin films.

Intensity dependence of photoconductivity has been studied at 303 K to investigate the nature of the recombination process in Se<sub>80-x</sub>Te<sub>20</sub>Bi<sub>x</sub> thin films. The plots of  $\ln \sigma_{ph}$  and  $\ln F$  are straight for all compositions as shown in Fig. 3, indicating that  $\sigma_{ph}$  follows a power law with  $F$ , i.e.  $\sigma_{ph} \propto F^\gamma$ , where  $\gamma$  is an exponent which characterizes the recombination mechanism. For  $\gamma = 0.5$ , recombination is bimolecular in nature, whereas for  $\gamma = 1$ , recombination is monomolecular [35]. In our case,  $\gamma$  varies from 0.51 to 0.56 for all compositions, which is close to 0.5. This confirms the dominance of bimolecular recombination mechanism and which may be discussed on the basis of Street and Mott model [35]. These authors suggest that chalcogenide glasses contains higher number of frozen-in structural defects at chalcogen atom sites, leaving broken or dangling bonds on the neighboring atoms which could be the source of localized defect states. The dangling bond

states are either unoccupied or occupied by pairs of electrons so that paired  $D^+$  and  $D^-$  defects are formed, which have lower energy than the neutral singly occupied sites, according to the relation:

$$2D^0 = D^+ + D^-$$

The energy level of the neutral unpaired dangling bond  $D^0$  lies intermediate between that of  $D^+$  and  $D^-$ . The  $D^+$  and  $D^-$  centres can act as discrete traps for electrons and holes generated by light excitation. This capture of photoexcited charge carriers leads to an excess concentration of  $D^0$  centres. These trapped carriers can be released to the valance or conduction band enhancing the conductivity.

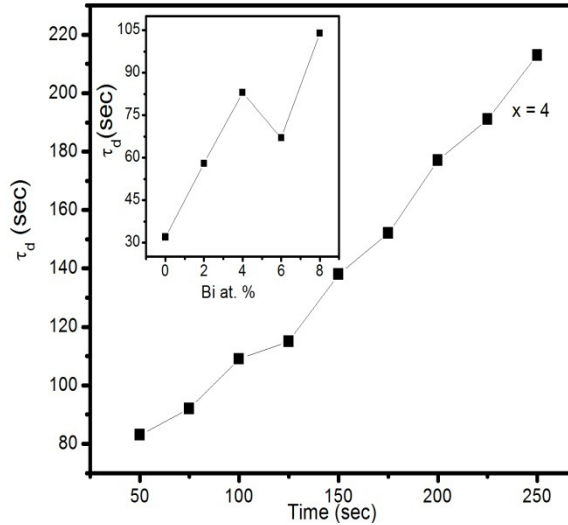


Fig. 4. Plot of  $\tau_d$  verses time for  $Se_{80-x}Te_{20}Bi_x$  thin films. Inset: Plot of  $\tau_d$  verses Bi at. % for  $Se_{80-x}Te_{20}Bi_x$  thin films.

Transient photoconductivity measurements were conducted by exposing all samples to a light intensity of 1043 lux at 303 K and amorphous thin films of a- $Se_{80-x}Te_{20}Bi_x$  show slow decay of photocurrent (results not shown here). A persistent photocurrent (the asymptotic value of the current in the decay) is also measured for all samples. The persistent photocurrent was subtracted from the measured photocurrent and found that the decay of photocurrent is non exponential even after subtracting the persistent photocurrent. In discussing the non-exponential decay process we prefer to use the concept of differential lifetime as described by Fuhs and Meyer [36]. According to these researchers, the differential life time can be defined as follows:

$$\tau_d = - \left[ \frac{1}{I_{ph}} \left[ \frac{dI_{ph}}{dt} \right] \right]^{-1} \quad (3)$$

In the case of an exponential decay, the differential life time will be equal to the carrier lifetime. However, in the case of a non-exponential decay,  $\tau_d$  will increase with time and only its value at  $t = 0$  corresponds to the carrier lifetime. The values of  $\tau_d$  are calculated at various times using slopes of  $I_{ph}$  verses time curves. Fig. 4 shows a plot of the variation of  $\tau_d$  with time at 303 K and 1043 lux intensity for  $x = 4$  at. % sample. It is clear from the figure that  $\tau_d$  increases with time. This confirms the non-exponential decay in the present case because for exponential decay  $\tau_d$  should be constant with time. Other compositions also follow a similar behaviour in their variation of  $\tau_d$  with time. To compare the value of  $\tau_d$  for the different samples, at 303 K the value of  $\tau_d$  at  $t = 50$  s verses the Bi content is plotted in the insert of Fig. 4. The value of  $\tau_d$  first increases as the concentration of Bi increases upto  $x = 4$  at. % and decreases from  $x = 4$  at. % to  $x = 6$  at. % and then again increases from  $x = 6$  at. % to  $x = 8$  at. %. The higher value of  $\tau_d$  after the Bi incorporation, indicating a slower rate of decay of  $I_{ph}$  and hence an increase in density of localized

states. One can therefore expect a higher density of localized states in the mobility gap of  $\text{Se}_{80-x}\text{Te}_{20}\text{Bi}_x$ .

#### 4. Conclusion

The effect of Bi additive on the dark ( $\sigma_d$ ) and photoconductivity ( $\sigma_{ph}$ ) is studied in amorphous  $\text{Se}_{80}\text{Te}_{20}$  binary alloy. The study of dark conductivity of  $\text{Se}_{80-x}\text{Te}_{20}\text{Bi}_x$  thin films as a function of temperature (273-343 K) reveals that conduction is an activated process with single activation energy. The charge carrier concentration ( $n_\sigma$ ) increases upto  $x = 4$  at. % and these carriers act as localized states due to which dark conductivity ( $\sigma_d$ ) and photoconductivity ( $\sigma_{ph}$ ) increases after the addition of Bi. The increase in dark conductivity is also interpreted in terms of increased hopping conduction due to increase in the concentration of localized states. The photosensitivity ( $\sigma_{ph}/\sigma_d$ ) decreases after the incorporation of Bi into the  $\text{Se}_{70}\text{Te}_{30}$  system, due to increase in the number of localized states. Intensity dependence photoconductivity has also been studied showing the existence of bimolecular recombination in the thin films. The transient photoconductivity measurements follow the non-exponential decay (for all samples) as differential life time ( $\tau_d$ ) increases with time. The differential life time also increases with increase of Bi content upto  $x = 4$  at. % indicating increase in density of localized states. As concentration of Bi is further increases (when  $x = 6$  at. %) a decrease in dark conductivity, photoconductivity, carrier concentration and an increase in dark activation energy, photosensitivity can be understood in term of structural changes because addition of Bi may satisfy the bonding states of Se-Te due to which density of localized states decreases and then with further increase of Bi content in  $\text{Se}_{80-x}\text{Te}_{20}\text{Bi}_x$  amorphous thin films (when  $x = 8$  at. %), the dark conductivity ( $\sigma_d$ ), charge carrier concentration ( $n_\sigma$ ) and photoconductivity ( $\sigma_{ph}$ ) again increase and the dark activation energy ( $\Delta E$ ) and photosensitivity ( $\sigma_{ph}/\sigma_d$ ) decrease.

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